New Results Involving a Class of Generalized Hurwitz-Lerch Zeta Functions and Their Applications

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Received September 15, 2013; Revised October 22, 2013; Accepted November 01, 2013

Abstract In this paper, we study a certain class of generalized Hurwitz-Lerch zeta functions. We derive several new and useful properties of these generalized Hurwitz-Lerch zeta functions such as (for example) their partial differential equations, new series and Mellin-Barnes type contour integral representations involving Fox’s $H$-function and a pair of summation formulas. More importantly, by considering their application in Number Theory, we construct a new continuous analogue of Lippert’s Hurwitz measure. Some statistical applications are also given.

Keywords: Hurwitz-Lerch zeta function, arithmetic density of number theory, partial differential equations, series and Mellin-Barnes type contour integral representations, Fox’s $H$-function, summation formulas, generalized Hurwitz measure, probability density function, moment generating function


1. Introduction, Definitions and Preliminaries

We begin by recalling the familiar general Hurwitz-Lerch Zeta function $\Phi(z,s,a)$, which is defined by (see, for example, [2]; see also [18,19,20,21]).

$$\Phi(z,s,a) := \sum_{n=0}^{\infty} \frac{z^n}{(n+a)^s} \quad (1.1)$$

($a \in \mathbb{C} \setminus \mathbb{Z}_0^0; s \in \mathbb{C}$, when $|z|<1$; $\Re(s)>1$ when $|z|=1$).

Special cases of the Hurwitz-Lerch Zeta function $\Phi(z,s,a)$ include (for example) the Riemann Zeta function $\zeta(s)$ and the Hurwitz (or generalized) Zeta function $\zeta(s,a)$ defined by (see, for details, [2, Chapter I] and [21, Chapter 2])

$$\zeta(s) := \sum_{n=1}^{\infty} \frac{1}{n^s} = \Phi(1,s,1) = \zeta(s,1) \quad (\Re(s)>1) \quad (1.2)$$

and

$$\zeta(s,a) := \sum_{n=0}^{\infty} \frac{1}{(n+a)^s} = \Phi(1,s,a) \quad (\Re(s)>1; a \in \mathbb{C} \setminus \mathbb{Z}_0^0), \quad (1.3)$$

respectively. Just as its aforementioned special cases $\zeta(s)$ and $\zeta(s,a)$, the Hurwitz-Lerch Zeta function $\Phi(z,s,a)$ defined by (1.4) can be continued meromorphically to the whole complex $s$-plane, except for a simple pole at $s = 1$ with its residue 1. It is also known that [[2], Equation 1.11 (3)]

$$\Phi(z,s,a) = \frac{1}{\Gamma(s)} \int_0^\infty e^{-t} e^{-at} \left(1-ze^{-t}\right)^{-s} dt$$

(1.4)

($\Re(s)>1; \Re(s)>0$, when $|z|\leq 1 (\neq 1); \Re(s)>1$ when $z=1$).

Recently, the following modified (and slightly generalized) version of the integral in (1.4) was introduced and studied by Raina and Chhajed [[16], Equation (1.6)]

$$\Phi^\mu(z,s,a,b) = \frac{1}{\Gamma(s)} \int_0^\infty e^{b} e^{-t} \left(1-ze^{-t}\right)^{-s} \mu dt \quad (\Re(s)>0; \Re(s)>0; \lambda > 0; \mu \in \mathbb{C}) \quad (1.5)$$

where we have assumed further that

$$\Re(s)>0 \text{ when } b=0 \text{ and } |z|\leq 1 \quad (\zeta \neq 1)$$

or

$$\Re(s-\mu)>0 \text{ when } b=0 \text{ and } \zeta = 1,$$
provided, of course, that the integral in (1.5) exists. As a matter of fact, the aforementioned investigation by Raina and Chhajed [16] was motivated by the following special case of the function \( \Theta^\mu_\mu(\zeta,s,a;b) \) defined by (1.5):

\[
\Theta^\mu_\mu(\zeta,s,a;0) = \Phi^\mu_\mu(\zeta,s,a) \nonumber \]

\[
= \frac{1}{\Gamma(s)} \int_0^\infty \frac{t^{s-1}e^{-at}}{(1-\zeta e^{-t})^\mu} \, dt \tag{1.6}
\]

( \( \Re(a) > 0; \Re(s) > 0 \) , when \( |\zeta| \leq 1(\zeta \neq 1) \); \( \Re(s-\mu) > 0 \) when \( \zeta = 1 \).) where the function \( \Phi^\mu_\mu(\zeta,s,a) \) defined by

\[
\Phi^\mu_\mu(\zeta,s,a) = \sum_{n=0}^\infty \frac{(\mu)_n \zeta^n}{(a+n)^\beta n!} \tag{1.7}
\]

was studied by Goyal and Laddha [14]. (Equation (1.5)).

Here, and in what follows, \( (\lambda)_n(\lambda, \nu \in \mathbb{C}) \) denotes the Pochhammer symbol (or the shifted factorial) which is defined, in terms of the familiar Gamma function, by

\[
(\lambda)_n = \frac{\Gamma(\lambda+n)}{\Gamma(\lambda)} \quad (n \in \mathbb{N}; \lambda \in \mathbb{C}),
\]

where it is understood conventionally that \( (0)_0 = 1 \) and assumed tacitly that the G-quotient exists, \( \mathbb{N} \) being the set of positive integers.

It may be of interest to observe in passing that, in terms of the Riemann-Liouville fractional derivative operator \( D^\mu_\zeta \) defined by (see, for example, [3,7,17])

\[
D^\mu_\zeta \left\{ f(\zeta) \right\} = \frac{1}{\Gamma(-\mu)} \int_0^\zeta (\zeta-t)^{-\mu-1} f(t) \, dt \quad (\Re(\mu) < 0)
\]

one may readily yield

\[
\Phi^\mu_\mu(\zeta,s,a) = \frac{1}{\Gamma(s)} \zeta^{s-1} \Phi(\zeta,s,a) \quad (\Re(\mu) > 0),
\]

which (as already remarked by Lin and Srivastava [8]) exhibits the interesting (and useful) fact that the function \( \Phi^\mu_\mu(\zeta,s,a) \) is essentially a Riemann-Liouville fractional derivative of the classical Hurwitz-Lerch function \( \Phi(\zeta,s,a) \).

One other special case of the function \( \Theta^\mu_\mu(\zeta,s,a;b) \) defined by (1.5) occurs when we set \( \lambda = \mu = 1 \) and \( \zeta = 1 \) in the definition (1.5). We thus obtain

\[
\Theta^1_1(1,s,a;b) = \zeta_b(s,a) \nonumber = \frac{1}{\Gamma(s)} \int_0^\infty \frac{e^{-at} - b}{1-e^t} \, dt. \tag{1.9}
\]

where \( \zeta_b(s,a) \) is the extended Hurwitz zeta function defined in [1]. In fact, just as it is already pointed out in [10], the series representation (see [16], Equation (2.1)) given for the function \( \Theta^\mu_\mu(\zeta,s,a;b) \) in (1.5) is incorrect. Obvious further specializations in (1.6) and (1.9) would immediately relate these functions with the Riemann zeta function \( \zeta(s) \) and the Hurwitz (or generalized) zeta function \( \zeta(s,a) \) defined by (1.2) and (1.3), respectively.

By using the series expansion of the binomial

\[
\left(1-\zeta e^{-t}\right)^{-\mu}
\]

occurring in the integrand of (1.5) and evaluating the resulting integral by means of the corrected version of a known integral formula [13], Equation (1.53) in terms of Fox’s \( H \)-function defined by (1.12) below, the following series and Mellin-Barnes type contour integral representations of the function \( \Theta^\mu_\mu(\zeta,s,a;b) \) defined by (1.5) were obtained in [10]:

\[
\Theta^\mu_\mu(\zeta,s,a;b) = \frac{1}{2\pi i} \frac{e^{\mu \zeta}}{\Gamma(s)\Gamma(\mu-s)} \int_{C} \frac{\Gamma(s)\Gamma(\mu-s)}{(a-s)^s} \left(\frac{1}{\lambda}\right)^\mu \, ds \quad (\lambda > 0)
\]

and

\[
\Theta^\mu_\mu(\zeta,s,a;b) = \frac{1}{2\pi i} \frac{e^{\mu \zeta}}{\Gamma(s)\Gamma(\mu-s)} \int_{C} \frac{\Gamma(s)\Gamma(\mu-s)}{(a-s)^s} \left(\frac{1}{\lambda}\right)^\mu \, ds \quad (\lambda > 0)
\]

it being assumed that each member of the assertions (1.10) and (1.11) exists.

**Remark 1.** The \( H \)-functions involved on the right-hand sides of (1.10) and (1.11) are particular cases of the celebrated Fox’s \( H \)-function which is defined as follows.

**Definition 1.** The well-known Fox’s \( H \)-function is defined here by (see, for details, [13], Definition 1.1; see also [6,11,23,24])

\[
H_{p,q}^{m,n}(\zeta) = H_{p,q}^{m,n}(\zeta) \left[ \begin{array}{c} a_p, A_p \\ b_p, B_q \end{array} \right] 
\]

\[
H_{p,q}^{m,n}(\zeta) = H_{p,q}^{m,n}(\zeta) \left[ \begin{array}{c} a_n, A_n \\ b_n, B_q \end{array} \right] 
\]

\[
= \frac{1}{2\pi i} \int_L \Xi(s) \zeta^{-s} \, ds 
\]

\[
= \frac{1}{2\pi i} \int_L \Xi(s) \zeta^{-s} \, ds 
\]
where
\[ \Xi(\zeta) = \prod_{j=1}^{m} \Gamma(b_j + B_j \zeta) \prod_{j=1}^{n} \Gamma(1-a_j - A_j \zeta) \prod_{j=m+1}^{q} \Gamma(1-b_j - B_j \zeta) \prod_{j=n+1}^{p} \Gamma(a_j + A_j \zeta) \] (1.13)

Here \( \zeta \in \mathbb{C}\setminus\{0\} \) with \( \arg(\zeta) < \pi \)

In this section, we will show that the generalized Hurwitz-Lerch zeta function \( \Theta^\lambda_{\mu+1}(\zeta, s, a; b) \) defined by (1.5). We first derive a partial differential equation satisfied by the function in (1.11). We then obtain another series representation and related results for this generalized Hurwitz-Lerch zeta function. The results derived here are also applied in our investigations concerning the generalized Hurwitz-Lerch zeta measure and its related statistical concepts.

2. Differential Equation of the Generalized Hurwitz-Lerch Zeta Function \( \Theta^\lambda_{\mu} \)

In this section, we will show that the generalized Hurwitz-Lerch zeta function \( \Theta^\lambda_{\mu}(\zeta, s, a; b) \) satisfies a partial differential equation when the parameter \( \lambda \) is given by
\[ \lambda = \frac{1}{m} \quad (m \in \mathbb{N}). \]

We first prove the following lemma which will be used in the proof of our main theorem.

Lemma (Derivative Property). The following derivative formulas hold true:
\[ \Theta^\lambda_{\mu+1}(\zeta, s, a; b) = \frac{1}{\mu} \zeta d\zeta \Theta^\lambda_{\mu}(\zeta, s, a; b) \quad (\lambda > 0) \quad (2.1) \]

and
\[ \frac{1}{\mu} \zeta \zeta d\zeta \Theta^\lambda_{\mu+1}(\zeta, s, a; b; m) = \frac{1}{\mu} \zeta d\zeta \left( \frac{1}{\mu} \zeta \zeta d\zeta \Theta^\lambda_{\mu}(\zeta, s, a; b; m) \right) \quad (m \in \mathbb{N}). \quad (2.2) \]

Proof. The proofs of the derivative formulas (2.1) and (2.2) are direct. For example, by applying the series representation (1.10), one easily finds that
\[ \frac{1}{\mu} \zeta d\zeta \Theta^\lambda_{\mu}(\zeta, s, a; b) = \frac{1}{\lambda d\zeta} \sum_{n=1}^{\infty} \left( \frac{\mu}{a+n} \right)^n \left( \frac{1}{\lambda} \right)^{n-1} \frac{1}{n!} = \Theta^\lambda_{\mu+1}(\zeta, s, a; b) \quad (\lambda > 0), \]

which is precisely the first result (2.1) asserted by the Lemma. The second assertion (2.2) follows immediately from (2.1) upon setting
\[ \lambda = \frac{1}{m} \quad (m \in \mathbb{N}) \quad \text{and} \quad b \rightarrow b^m \quad (m \in \mathbb{N}). \]

Our first main result is contained in the following theorem.

Theorem 1. The generalized Hurwitz-Lerch zeta function \( \Theta^m_{\mu}(\zeta, s, a; b) \) \( (m \in \mathbb{N}) \) satisfies the following partial differential equation:
\[ \left[ (-1)^{m+1} m^m \mathcal{D}_b \right] \frac{1}{\mu} \zeta d\zeta \Theta^m_{\mu}(\zeta, s, a; b) = 0, \quad (2.3) \]

where the differential operators \( \mathcal{D}_b \) and \( \partial_\zeta \) are given by
\[ \mathcal{D}_b := \partial_b \left( \partial_b - s \right) \left( \partial_b - \frac{1}{m} \right) \cdots \left( \partial_b - \frac{m-1}{m} \right) \quad (2.4) \]

and
\[ \partial_\zeta := \zeta \partial_\zeta, \]

respectively.

Proof. The proofs of the derivative formulas (2.1) and (2.2) are direct. For example, by applying the series representation (1.10), one easily finds that
\[ \frac{1}{\mu} \zeta d\zeta \Theta^\lambda_{\mu}(\zeta, s, a; b) = \frac{1}{\lambda d\zeta} \sum_{n=1}^{\infty} \left( \frac{\mu}{a+n} \right)^n \left( \frac{1}{\lambda} \right)^{n-1} \frac{1}{n!} = \Theta^\lambda_{\mu+1}(\zeta, s, a; b) \quad (\lambda > 0), \]

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and
\[ \partial_\zeta := \zeta \partial_\zeta, \]

respectively.

Proof. The proofs of the derivative formulas (2.1) and (2.2) are direct. For example, by applying the series representation (1.10), one easily finds that
\[
\frac{1}{\lambda} = m \quad (m \in \mathbb{N}) \quad \text{and} \quad b \to b^m \quad (m \in \mathbb{N})
\]
in the above equation (2.5) and then applying the well-known (Gauss-Legendre) multiplication formula:
\[
\Gamma(mz) = (2\pi)^{1-m}m^{-\frac{1}{2}}\prod_{j=1}^{m-1}(\zeta + j - \frac{1}{m})
\]
(2.6)
we find that
\[
H^n_{2m}[(a-s)^b, (0, m)] = \frac{1}{\lambda^m} \int_0^1 \Gamma(s + w)\Gamma(mw)[(a-s)^b]^{-w} dw
\]
(2.7)
and
\[
\mathcal{G}^{m+1,0}_{0,m+1}\left((a-s)bm^{-m}\right)
\]
is a very specialized case of Meijer’s G-function \(G^{m,n}_{p,q}(z)\) which, in turn, is a special case of Fox’s H-function defined by (1.12), that is, we have the following relationship (see, for details, [14]; see also [2, 12, 15]):
\[
G^{m,n}_{p,q}(z) = G^{m,n}_{p,q}\left(\zeta \left(\begin{array}{c}
(a_j)_{j=1}^p \\
(b_j)_{j=1}^q
\end{array}\right)\right) = H^{m,n}_{p,q}(z)\left(\begin{array}{c}
(a_1, \ldots, a_p, 0, 0, \ldots, b_1, \ldots, b_q, 0, 0, \ldots)
\end{array}\right)
\]
(2.8)
We know that the function \(W\) defined by
\[
W := G^{m,n}_{p,q}(z)\left(\begin{array}{c}
(a_j)_{j=1}^p \\
(b_j)_{j=1}^q
\end{array}\right)
\]
satisfies the following differential equation of order \(\max(p, q)\) (see, for example, [12], Equation 5.4(1)):
\[
\left(-1\right)^{p-m+n}\left(\frac{\partial}{\partial \zeta} - a_1 + 1\right)\cdots\left(\frac{\partial}{\partial \zeta} - a_p + 1\right)
\]
\[
-\left(\frac{\partial}{\partial \zeta} - b_1\right)\cdots\left(\frac{\partial}{\partial \zeta} - b_q\right)
\]
\[W = 0,
\]
where
\[
\frac{\partial}{\partial \zeta} = \zeta \frac{d}{d\zeta}.
\]
Hence, clearly, the function given by (2.7) satisfies the following differential equation:
\[
\left(-1\right)^{m+1}(a-s)m^{-m}b
\]
\[
-\theta_b\left(\theta_b - s\right)\cdots\left(\theta_b - \frac{m-1}{m}\right)
\]
\[
\mathcal{G}^{m+1,0}_{0,m+1}\left((a-s)bm^{-m}\right)
\]
(2.9)
where, as already stated in Theorem 1,
\[
\theta_b = b \frac{\partial}{\partial b}
\]
Now, if we write (see also Equation (2.4))
\[
\mathcal{D}_b := \theta_b\left(\theta_b - \frac{1}{m}\right)\cdots\left(\theta_b - \frac{m-1}{m}\right)
\]
(2.10)
then the equation (2.9) becomes
\[
\mathcal{D}_b\left(\mathcal{G}^{m+1,0}_{0,m+1}\left((a-s)bm^{-m}\right)\right) = \left(-1\right)^{m+1}(a-s)m^{-m}b
\]
(2.10)
Applying the differential operator \(\mathcal{D}_b\) to the function \(\frac{1}{\lambda^{m+1}}\left(-\zeta, s, a; b^m\right)\) given by (1.11) with
\[
\zeta \to -\zeta, \lambda = \frac{1}{m} (m \in \mathbb{N}) \quad \text{and} \quad b \to b^m \quad (m \in \mathbb{N}),
\]
we find by using (2.10) that
\[
\mathcal{D}_b\left(\frac{1}{\lambda^{m+1}}\left(-\zeta, s, a; b^m\right)\right)
\]
\[
= \sqrt{m}\left(2\pi\right)^{1-m}m^{-\frac{1}{2}}\int_0^1 \Gamma(s)\Gamma(\mu - s)
\]
\[
\left(a - s\right)^{\zeta - \mu} d\zeta
\]
\[
= \frac{1}{\lambda^{m+1}}\left(-\zeta, s, a; b^m\right)
\]
(2.11)
where the first integral \(\mathcal{I}_1\) is actually the generalized Hurwitz-Lerch zeta function given by
\[
\mathcal{I}_1 = \frac{1}{\lambda^{m+1}}\left(-\zeta, s, a; b^m\right).
\]
(2.12)
The evaluation of the second integral \(\mathcal{I}_2\) given by
\[ J_2 = \frac{\sqrt{m} (2\pi)^{1-m}}{2\pi i} \int_{C} \frac{1}{\Gamma(s)} \frac{\Gamma(s+1) \Gamma(\mu-s)}{(a-s)^{\alpha}} \] (2.13)

is more complicated. Since the residues of \( \Gamma(s+1) \) at the poles \( s = -k \) \((k \in \mathbb{N})\) are computed by

\[ \text{Res}_{s=-k} \{ \Gamma(s+1) \} = \lim_{s \to -k} (s+k) \Gamma(s+1) = \frac{(-1)^{k-1}}{(k-1)!}, \] (2.14)

the Residue Theorem implies that

\[ J_2 = \frac{\sqrt{m} (2\pi)^{1-m}}{2\pi i} \int_{C} \frac{1}{\Gamma(s)} \frac{\Gamma(\mu - s)}{(a-s)^{\alpha}} \] (2.15)

Finally, by setting

\[ b \to b^m \quad (m \in \mathbb{N}) \quad \text{and} \quad \zeta \to -\zeta \]

in the last equation (2.18), we arrive at the desired result (2.3) asserted by Theorem 1.

It is interesting to consider a special case of Theorem 1 when \( m = 1 \). Thus, if we write

\[ \Theta^1_{\mu}(\zeta, s, a; b) = \Theta_{\mu}(\zeta, s, a; b), \] (2.19)

then we have the following corollary.

**Corollary 1.** The generalized Hurwitz-Lerch zeta function 

\[ \Theta_{\mu}(\zeta, s, a; b) \]

satisfies the following partial differential equation:

\[ b \frac{\partial}{\partial b} \left( b \frac{\partial}{\partial b} - s \right) \left( b \frac{\partial}{\partial b} - 1 \right) \] \( \Theta_{\mu}(\zeta, s, a; b) = 0 \). (2.20)

Furthermore, the function \( \Theta_{\mu}(1, s, a; b) \), considered as an analytic function of the variable \( b \) satisfies the following relation:

\[ b \frac{\partial}{\partial b} \left( b \frac{\partial}{\partial b} - s \right) \left( b \frac{\partial}{\partial b} - 1 \right) \Theta_{\mu}(1, s, a; b) = (a+1) b \mu \Theta_{\mu+1}(1, s, a+1; b) \] (2.21)

### 3. Further Series Representations and Related Results

In this section, we first give a new series representation of the generalized Hurwitz-Lerch zeta function \( \Theta^1_{\mu}(\zeta, s, a; b) \) involving the familiar Laguerre polynomials of order (index) \( \alpha \) and degree \( n \) in \( x \), which are generated by

\[ (1-t)^{-\alpha-1} \exp \left( -\frac{xt}{1-t} \right) = \sum_{n=0}^{\infty} L_n^{(\alpha)}(x) t^n \quad (|t| < 1; \alpha \in \mathbb{C}). \] (3.1)
Indeed, upon setting $t = 1 - t^x$ and $x = b$
in (3.1), we get
\[
\exp \left( - \frac{b}{t^x} \right) = t^x \frac{d^b}{d b} \sum_{n=0}^{\infty} \frac{x}{t^n} \frac{d}{dx} \left( 1 - t^x \right)^n. \tag{3.2}
\]

We now make use of (3.2) and the series expansions of the binomials
\[
\left(1-t^x\right)^n \quad \text{and} \quad (1-\varepsilon t^{-1})^{-\mu}
\]
occuring in the integrand of (1.5). By evaluating the resulting Eulerian integral, we thus arrive at the series representations given by Theorem 2 below.

**Theorem 2.** Each of the following series representations holds true for the generalized Hurwitz-Lerch zeta function $\Theta_\mu(z, s; a; b)$:
\[
\Theta_\mu(z, s; a; b) = \frac{e^{-b}}{\Gamma(s)} \sum_{n=0}^{\infty} \sum_{j=0}^{\infty} \frac{(n)!}{n!} \left( \frac{\mu + t - 1}{t} \right)^n \Gamma(s + \lambda(\alpha + j + 1)) \cdot \Phi_\mu(z, s + \lambda(\alpha + j + 1), a) \quad (\mathcal{R}(a) > 0; \mathcal{R}(s + \lambda(\alpha + j + 1)) > -\lambda), \tag{3.3}
\]
and
\[
\Theta_\mu^*(z, s; a; b) = \frac{e^{-b}}{\Gamma(s)} \sum_{n=0}^{\infty} \sum_{j=0}^{\infty} \frac{(n)!}{n!} \left( \frac{\mu + t - 1}{t} \right)^n \Gamma(s + \lambda(\alpha + j + 1)) \cdot \Phi_\mu^*(z, s + \lambda(\alpha + j + 1), a) \quad (\mathcal{R}(a) > 0; \mathcal{R}(s + \lambda(\alpha + j + 1)) > -\lambda), \tag{3.4}
\]
provided that each member of the assertions (3.3) and (3.4) exists, $\Theta_\mu^*(z, s; a; b)$ being given by (1.7).

**Proof.** As already outlined above, our demonstration of the first assertion (3.3) of Theorem 2 is based essentially upon the representation (3.2) and the following well-known Eulerian integral:
\[
\int_0^\infty t^{\mu-1} e^{-\sigma t} dt = \left( \frac{\Gamma(\rho)}{\sigma^\rho} \right) \left( \min \{ \mathcal{R}(\rho), \mathcal{R}(\sigma) \} > 0 \right). \tag{3.5}
\]

The second assertion (3.4) follows from the first assertion (3.3) when we interpret the $t$-series in (3.4) by means of the definition (1.7).

In our derivation of each of the summation formulas (3.3) and (3.4), it is assumed that the required inversions of the order of summation and integration are justified by absolute and uniform convergence of the series and integrals involved. The final results (3.3) and (3.4) would thus hold true whenever each member of the assertions (3.3) and (3.4) of Theorem 2 exists. □

**Remark 2.** For the extended Hurwitz zeta function $\zeta_\mu(z, s; a)$ defined by (1.9), it is easily deduced from the assertion (3.4) of Theorem 2 when $\lambda = \mu = 1$ and $z = 1$ that
\[
\zeta_\mu(z, s; a) = \frac{e^{-b}}{\Gamma(s)} \sum_{n=0}^{\infty} \sum_{j=0}^{\infty} \frac{(-1)^j}{(n+j)!} \Gamma(s + \alpha + j + 1) \cdot \Phi_\mu(z, s + \alpha + j + 1, a) \quad (s > 0; \mathcal{R}(s + \alpha + j + 1) > 0), \tag{3.6}
\]
provided that each member of (3.6) exists, $\zeta_\mu(z, s; a)$ being the Hurwitz (or generalized) zeta function given by (1.3). The obvious further special case of (3.6) when $a = 1$ and $\alpha = 0$ would yield the corrected version of a known result (see [11], Equation (7.78)).

We now give a pair of summation formulas involving the generalized Hurwitz-Lerch zeta function $\Theta_\mu^{\pm}(z, s; a; b)$.

**Theorem 3.** Each of the following summation formulas holds true for the generalized Hurwitz-Lerch zeta function $\Theta_\mu^{\pm}(z, s; a; b)$:
\[
2^{s-1} \left[ \Theta_\mu^{\pm}(z, s; a; b) + \Theta_\mu^{\pm}(z, s; a; b) \right] = \sum_{n=0}^{\infty} \frac{(-\mu)^n}{n!} \Theta_\mu^{\pm}(z, s + a, b + n; 2^2 b) \zeta^{2n} \tag{3.7}
\]
and
\[
2^{s-1} \left[ \Theta_\mu^{\pm}(z, s; a; b) - \Theta_\mu^{\pm}(z, s; a; b) \right] = \sum_{n=0}^{\infty} \frac{(-\mu)^{2n+1}}{(2n+1)!} \Theta_\mu^{\pm}(z, s + a + 1, b + n; 2^2 b) \zeta^{2n+1}, \tag{3.8}
\]
provided that each member of the assertions (3.7) and (3.8) exists.

**Proof.** Making use of the integral representation in (1.5) for the function $\Theta_\mu^{\pm}(z, s; a; b)$, we get
\[
\Theta_\mu^{\pm}(z, s; a; b) + \Theta_\mu^{\pm}(z, s; a; b) = \frac{1}{\Gamma(s)} \int_0^\infty t^{\mu-1} \exp \left( -a t - b \frac{dt}{t} \right) \left( 1 - \varepsilon t^{-1} \right)^\mu + \left( 1 + \varepsilon t^{1-\mu} \right)^\mu \tag{3.9}
\]
Since
\[
\left( 1 - \varepsilon t^{-1} \right)^\mu + \left( 1 + \varepsilon t^{1-\mu} \right)^\mu = 2 \sum_{n=0}^{\infty} \frac{(-\mu)^n}{n!} \zeta^{2n} e^{-2nt}, \tag{3.10}
\]
by substituting from (3.10) into (3.9) and interchanging the order of summation and integration, we find that
\[
\Theta_\mu^{\pm}(z, s; a; b) + \Theta_\mu^{\pm}(z, s; a; b) = \frac{2}{\Gamma(s)} \sum_{n=0}^{\infty} \frac{(-\mu)^n}{n!} \zeta^{2n} \int_0^\infty t^{\mu-1} \exp \left( -a t^2 - b \frac{dt}{t^2} \right) \left( 1 - \zeta e^{-2t} \right)^\mu dt
\]
\[
= \frac{2^{s-1}}{\Gamma(s)} \sum_{n=0}^{\infty} \frac{(-\mu)^n}{(2n)!} \zeta^{2n} \exp \left( -a t^2 - b \frac{dt}{t^2} \right) \left( 1 - \zeta e^{-2t} \right)^\mu dt
\]
\[
= \frac{2^{s-1}}{\Gamma(s)} \sum_{n=0}^{\infty} \frac{(-\mu)^n}{(2n)!} \Theta_\mu^{\pm}(z, s + a, b + n; 2^2 b) \zeta^{2n}, \tag{3.11}
\]
which obviously proves the assertion (3.7) of Theorem 3.

The assertion (3.8) of Theorem 3 can be proven in a manner analogous to that detailed above. □
Remark 3. If we set $\mu = 1$ in (3.7) and (3.8), the series occurring on their right-hand sides would terminate. Upon setting $r = -s$, we thus obtain

$$2^{s-1} \left[ \Theta^s (r, s; 2a; b) + \Theta^s (r, s; 2a; b) \right]$$

and

$$2^{s-1} \left[ \Theta^s (r, s; 2a; b) - \Theta^s (r, s; 2a; b) \right]$$

In particular, if we set $\zeta = 1$ in these last two summation formulas (3.12) and (3.13), we get

$$\Theta^s _1 (1, s, a; 2^2 b) = \Theta^s _1 (1, s, a; 2^2 b) - \Theta^s _1 (1, s, 2a; b)$$

and

$$\Theta^s _2 (1, s, 2a; b) = \Theta^s _2 (1, s, 2a; b) - \Theta^s _1 (1, s, a; 2^2 b)$$

respectively. In its further special case when $\lambda = 1$, the summation formula (3.14) can be shown to correspond to known results (see, for example, [11], Theorem 7.9.; see also [1]).

4. A Generalization of the Hurwitz Measure

Suppose that $\chi_A(n)$ denote the characteristic function of the subset $A$ of the set $\mathbb{N}$ of positive integers (or, in the language of probability theory, the indicator function of the event $A \subseteq \mathbb{N}$). Then it is well known that the following arithmetic density of number theory:

$$\text{dens}(A) = \lim _{k \to \infty} \frac{1}{k} \sum _{n=1} ^k \chi_A(n)$$

does not define a measure on the set $\mathbb{N}$ of positive integers. In order to remedy this deficiency, Golomb [5] defined a probability on the sample space $\mathbb{N}$ as follows:

$$Q_s (A) = \frac{1}{\zeta(s)} \sum _{n=1} ^N \chi_A(n)$$

where $\zeta(s)$ denotes the Riemann zeta function defined by (1.2) and the characteristic (or indicator) function $\chi_A(n)$ is given by

$$\chi_A(n) = \begin{cases} 1 & (n \in A) \\ 0 & (n \not\in A) \end{cases}$$

Furthermore, Golomb [5] showed that, if the subset $A$ of $\mathbb{N}$ has an arithmetic density, then

$$\lim _{s \to 1^+} Q_s (A) = \text{dens}(A)$$

by allowing number-theoretic facts regarding densities of sets of positive integers to be proven by probabilistic means and then showing that such properties are preserved in the limit.

In an interesting sequel to Golomb’s investigation [5], Lippert [9] gave an analogous definition of the probabilities $P_s$ when the set $\mathbb{N}$ is replaced by the set of all real numbers greater than 1. Thus, for a Borel set $A \subseteq (1, \infty)$, Lippert’s Hurwitz measure of the set $A$ is defined by (see, for details, [9], Definition 1)

$$P_s (A) = \int _{\chi \in (1, \infty)} \chi_A(\chi) \zeta(s+1, \chi) d\chi$$

or, equivalently, by

$$P_s (A) = \int _{x \in (1, \infty)} \chi_A(x) d\tilde{\mu}(x, s)$$

where, in terms of the Hurwitz (or generalized) zeta function $\zeta(s, a)$ defined by (1.3),

$$\tilde{\mu}(x, s) := \frac{\zeta(s, x)}{\zeta(s)}$$

and

$$d\tilde{\mu}(x, s) = d\zeta (\mu, s) = d\frac{\zeta(s+1, x)}{\zeta(s)}$$

In this section, we propose to introduce a new continuous analogue of Lippert’s Hurwitz measure in (4.5) by using a special case of the generalized Hurwitz-Lerch zeta function $\Theta^s_\mu (\zeta, s, a; b)$ defined by (1.5).

Definition 2. For a Borel set $A \subseteq (1, \infty)$, the generalized Hurwitz measure of the set $A$ is defined by

$$P_s (A) = \int _{\chi \in (1, \infty)} \chi_A(\chi) \Theta^s _1 (1, s+1, a; b) d\chi$$

or, equivalently, by

$$P_s (A) = \int _{x \in (1, \infty)} \chi_A(x) d\tilde{\mu}(x, s; a, \lambda)$$

where

$$\tilde{\mu}(x, s; a, \lambda) := \frac{\Theta^s _2 (1, s, a; b)}{\Theta^s _1 (1, s, 1; b)}$$

and

$$d\tilde{\mu}(x, s; a, \lambda) = d\frac{\Theta^s _2 (1, s, a; b)}{\Theta^s _1 (1, s, 1; b)}$$

since it is easily seen from the definition (1.5) that

$$d\frac{\Theta^s _1 (\zeta, s, a; b)}{\Theta^s _1 (\zeta, s+1, a; b)} = -s \Theta^s _1 (\zeta, s+1, a; b)$$
Theorem 4. If \( A = (a, a+1) \), then
\[
P_s(A) = P_s((a, a+1))
\]
\[
H_{0,2}^{2.0} \left[ \left( \frac{1}{ab^2} \right) \left( s, 1, \left( 0, \frac{1}{a} \right) \right) \right] = \frac{\lambda a^2 \Gamma(s) \Theta_0^2(1, s, 1; b)}{\Theta_0^2(1, s, 1; b)}
\]
More generally, the generalized Hurwitz measure of an open set \( A \subseteq (1, \infty) \) is given by
\[
P_s(A) = \sum_{i \in I} P_s((a_i, b_i))
\]
where \( A = \bigcup_{i \in I} (a_i, b_i) \) \( \bigcup \bigcup (a_i, b_i) \in [1, \infty); i \in I \).

Theorem 5. Corresponding to the generalized Hurwitz measure given by (4.19), let
\[
\lambda \Gamma(s) = \sum_{i \in I} \Theta_0^2(1, s, 1; b) - \Theta_0^2(1, s, 1; b + \varepsilon)
\]

Proof. From (4.19), we have
\[
P_s(A) = \sum_{i \in I} \left( \Theta_0^2(1, s, 1; b) - \Theta_0^2(1, s, 1; b + \varepsilon) \right)
\]

By expanding the function \( \Theta_0^2(1, s, 1; b) \) by means of Taylor’s series and using the derivative formula (4.12), we get
\[
P_s(A) = \sum_{i \in I} \left( \Theta_0^2(1, s, 1; b) - \frac{\varepsilon^2}{2} s (s+1) \sum_{i=1}^{\infty} \Theta_0^2(1, s+2, i; b) + \cdots \right)
\]

We now consider each sum in (4.23) separately. We thus find that
\[
\sum_{i=1}^{\infty} \Theta_0^2(1, s+m, i; b)
\]

Remark 5. For \( \lambda = 1 \) and by letting \( b \to 0 \), the difference equation (4.15) reduces to the following form:
\[
\dot{\mu}(a+1, s) - \dot{\mu}(a, s) = \frac{1}{a^2} \zeta(s)
\]

where \( \dot{\mu}(x, s) \) is given by (4.7).

For open events, the generalized Hurwitz measure \( P_s(A) \) in (4.8) or (4.9) can be evaluated by using (4.9) and the above Proposition. The results are being stated as Theorem 4 below.
Since the number of non-negative integer solutions of the Diophantine equation \( j + n = N \) is
\[
\binom{N + 1}{1} = N + 1,
\]
the double summation in (4.24) can be replaced by a single summation, that is,\[
\sum_{i=1}^{\infty} \Theta^d_i (1,s + m,i;b) = \frac{1}{\lambda} \sum_{N=0}^{\infty} H^{2,0}_{0,2} \left( (N+1)b, s+m,1,0,1 \right)
\]
\[
= \Theta^d_i (1,s + m-1,1,b).
\]
(4.25)
We thus obtain\[
\lim_{s \to 1} \frac{P_s(A)}{s} = \lim_{s \to 1} \frac{\Theta^d_i (1,s+1,b)}{\Theta^d_i (1,s+1,b) + \cdots} \]
\[
= \epsilon - \frac{\epsilon^2}{2(s+1)} \lim_{s \to 1} \Theta^d_i (1,s+1,b) + \cdots.
\]
(4.26)
We note that, when \( s \to 1 \), the series for \( \Theta^d_i (1,s+1,b) \) is divergent and the series for \( \Theta^d_i (1,s+1,b) \) is convergent. Therefore, all other terms vanish in (4.26) except the leading term. Consequently, we get\[
\lim_{s \to 1} P_s(A) = \epsilon,
\]
(4.27)
which completes the proof of Theorem 5. \( \square \)

As in the theory of probability, we introduce the following definition.

**Definition 3.** A random variable \( \xi \) is said to be generalized Hurwitz distributed if its probability density function (p.d.f.) is given by\[
f_\xi (a) = \begin{cases} s \Theta^d_i (1,s+1,a;b) / \Theta^d_i (1,s+1,b) & (a \geq 1) \\ 0 & (\text{otherwise}) \end{cases}
\]
(4.28)

**Theorem 6.** Let \( \xi \) be a continuous random variable \( \xi \) with its p.d.f defined by (4.28). Then the moment generating function \( M(\xi) \) of the random variable \( \xi \) is given by\[
M(\xi) = E_s \left[ e^{\xi^a} \right] = \sum_{n=0}^{\infty} E_s \left[ \frac{\xi^n}{n!} \right]
\]
(4.29)
with the moment \( E_s \left[ \xi^n \right] \) of order \( n \) given by\[
E_s \left[ \xi^n \right] = n! \sum_{k=0}^{n} \frac{\Gamma(s-k) \Theta^d_i (1,s-k,1;b)}{\Gamma(s) \Theta^d_i (1,s,1;b)}.
\]
(4.30)

**Proof.** The assertion in (4.29) follows easily by using the exponential series for \( e^{\xi^a} \). If we use integration by parts, we find from the definition that

\[ E_s \left[ \xi^n \right] = \int_0^\infty a^n f_\xi (a) \, da \]
\[ = \frac{s}{\Theta^d_i (1,s+1,b)} \int_0^\infty a^n \Theta^d_i (1,s+1,a;b) \, da \]
\[ = \frac{1}{\Theta^d_i (1,s+1,b)} \int_0^\infty a^nd \left( \Theta^d_i (1,s,a;b) \right) \]
\[ = \frac{a^n \Theta^d_i (1,s,a;b)}{\Theta^d_i (1,s+1,b)} \int_0^\infty a^n \Theta^d_i (1,s+1,b) \, da \]
\[ \lim_{a \to \infty} \frac{n}{\Theta^d_i (1,s+1,b)} + 1 \]
\[ + \frac{n}{\Theta^d_i (1,s+1,b)} \int_0^\infty a^n \Theta^d_i (1,s,a;b) \, da \]
\[ = 1 + \frac{n}{\Theta^d_i (1,s+1,b)} \int_0^\infty a^n \Theta^d_i (1,s,a;b) \, da \quad (n \in \mathbb{N}),
\]
(4.31)
where we have also used the derivative property (4.12) and the following limit formula:
\[ \lim_{a \to \infty} \frac{a^n}{\Theta^d_i (1,s+1,b)} = 0, \quad (n \in \mathbb{N}).
\]

Consequently, we have the following reduction formula for \( E_s \left[ \xi^n \right] \):
\[ E_s \left[ \xi^n \right] = 1 + \frac{n}{\Theta^d_i (1,s+1,b)} \frac{n}{s-1} E_{s-1} \left[ \xi^{n-1} \right] \quad (n \in \mathbb{N}).
\]
(4.33)

By iterating the recurrence (4.33), we arrive at the desired result (4.30) asserted by Theorem 6. \( \square \)

**Remark 6.** The assertion (4.30) of Theorem 6 provides a generalization of a known result \([9],\) Proposition 3.

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**References**


